

Theory of gravitational waves

Laura Bernard

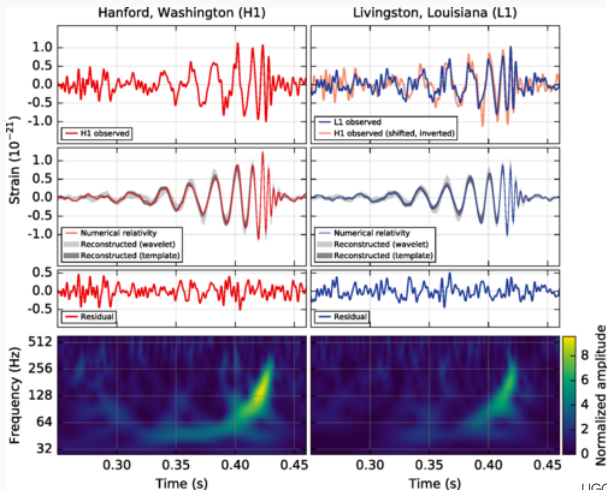
The first ACME workshop

The gravitational wave sky and complementary observations

7-11 April 2025, Toulouse



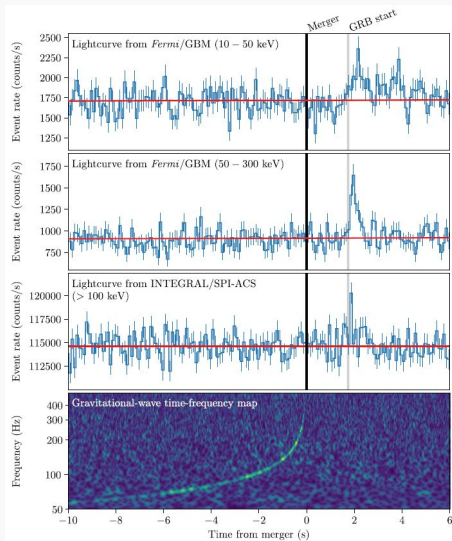
The first gravitational wave detections



LIGO-Virgo GW150914

$$\text{Chirp: } \frac{df}{dt} = \frac{96}{5} \pi^{8/3} \mathcal{M}^{5/3} f^{11/3}, \quad \mathcal{M} \equiv \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}}$$

The first gravitational wave detections



LIGO-Virgo GW170817

$$\text{Chirp: } \frac{df}{dt} = \frac{96}{5} \pi^{8/3} \mathcal{M}^{5/3} f^{11/3}, \quad \mathcal{M} \equiv \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}}$$

O4 and GW230529

LIGO/Virgo/KAGRA Public Alerts

- More details about public alerts are provided in the [LIGO/Virgo/KAGRA Alerts User Guide](#).
- Retractions are marked in **red**. Retraction means that the candidate was manually vetted and is no longer considered a candidate of interest.
- Less-significant events are marked in **grey**, and are not manually vetted. Consult the [LVK Alerts User Guide](#) for more information on significance in O4.
- Less-significant events are not shown by default. Press "Show All Public Events" to show significant and less-significant events.

O4 Significant Detection Candidates: **142** (158 Total - 16 Retracted)

O4 Low Significance Detection Candidates: **2484** (Total)

Show All Public Events

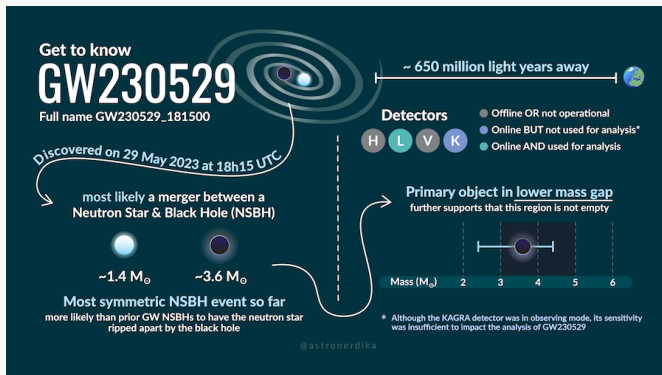
Page 1 of 11, next last »

SORT: EVENT ID (A-Z) ▼

..... ▼

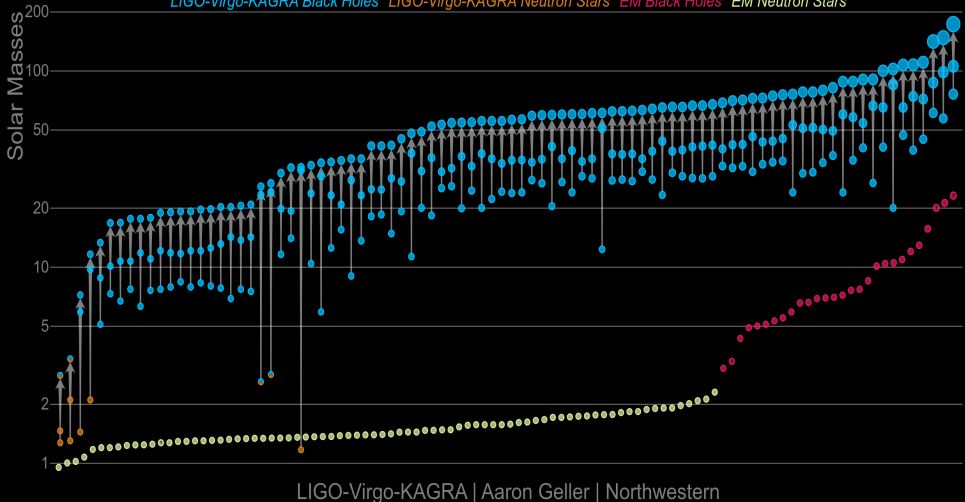
Event ID	Possible Source (Probability)	Significant	UTC	GCN	Location	FAR	Comments
S240930du	BBH (67%), Terrestrial (33%)	Yes	Sept. 30, 2024 23:46:14 UTC	GCN Circular Query Notices VOE		1 per 2.4747 years	
S240930aa	BBH (>99%)	Yes	Sept. 30, 2024 03:59:59 UTC	GCN Circular Query Notices VOE		1 per 1.0344e+11 years	
S240925n	BBH (>99%)	Yes	Sept. 25, 2024 00:58:09 UTC	GCN Circular Query Notices VOE		1 per 7.9146e+11 years	
S240924a	BBH (>99%)	Yes	Sept. 24, 2024 00:03:16 UTC	GCN Circular Query Notices VOE		1 per 12.869 years	
S240923ct	BBH (>99%)	Yes	Sept. 23, 2024 20:40:06 UTC	GCN Circular Query Notices VOE		1 per 4.1462e+07 years	
S240922df	BBH (>99%)	Yes	Sept. 22, 2024 14:21:06 UTC	GCN Circular Query Notices VOE		1 per 2.2729e+16 years	
S240921cw	BBH (>99%)	Yes	Sept. 21, 2024 20:18:35 UTC	GCN Circular Query Notices VOE		1 per 39.517 years	
S240920dw	BBH (>99%)	Yes	Sept. 20, 2024 00:00:00 UTC	GCN Circular Query Notices VOE		1 per 3.2668e+43 years	

O4 and GW230529



Masses in the Stellar Graveyard

LIGO-Virgo-KAGRA Black Holes *LIGO-Virgo-KAGRA Neutron Stars* *EM Black Holes* *EM Neutron Stars*



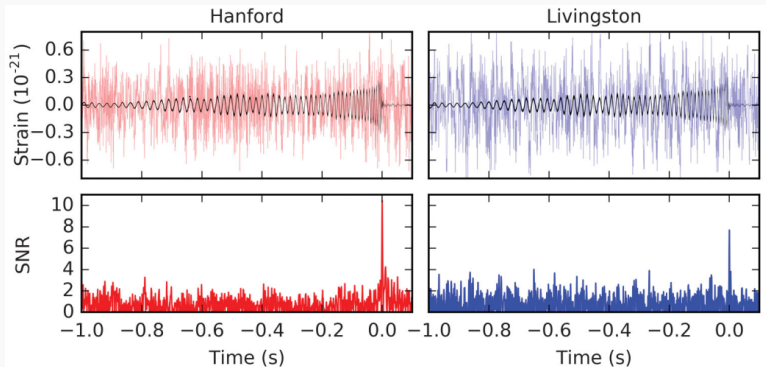
Why do we need...

1. ...to have a bank of **extremely precise** waveform templates?
2. ...to use **different modeling techniques**?
3. ...to go **beyond GR**?



Why do we need...

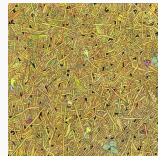
- to have a bank of **extremely precise** waveform templates?



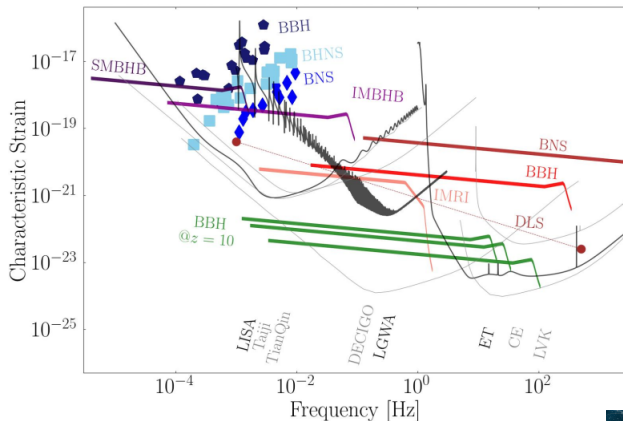
LIGO-Virgo GW151226

* chercher une aiguille dans une botte de foin

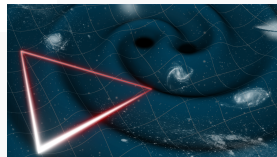
* looking for a needle in a haystack



The future gravitational wave universe



Einstein Telescope blue book (2025)

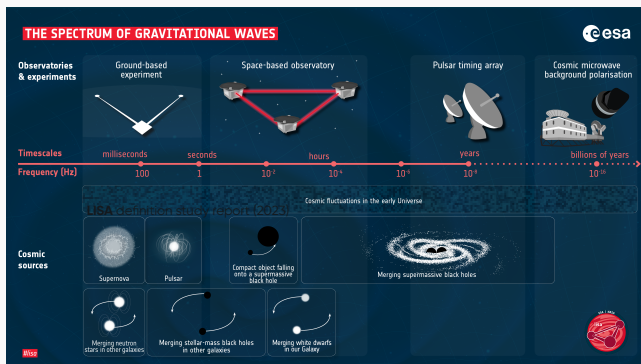


Why do we need...

1. ...to have a bank of **extremely precise** waveform templates?
 - ▶ detection and parameter estimation
2. ... to use **different modeling techniques**?

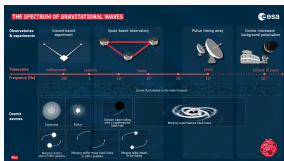
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 - detection and parameter estimation
2. ... to use **different modeling techniques**?
 - very different sources

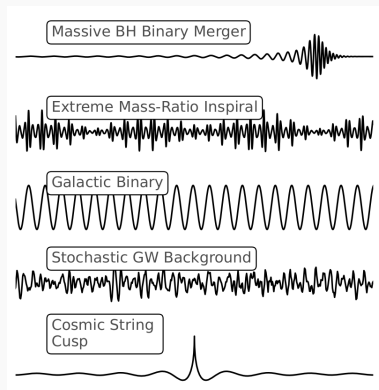


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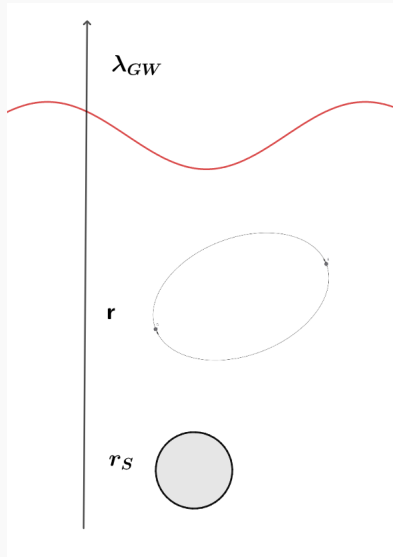


LISA definition study report (2023)



- GR highly non linear $\Rightarrow$ need **numerical** and **analytical** calculations

The b.a-ba of gravitational waves



The Einstein field equations

$$R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R = 8\pi T^{\mu\nu}$$

The hierarchy of scales

$$r_S \ll r \sim \frac{r_S}{v} \ll \lambda_{GW} \sim \frac{r}{v}$$

The b.a-ba of gravitational waves

$$h^{\mu\nu} \equiv \sqrt{-g} g^{\mu\nu} - \eta^{\mu\nu}$$

$$\underbrace{\square h^{\mu\nu}}_{\text{flat d'Alembertian } \eta^{\rho\sigma} \partial_\rho \partial_\sigma h} = \overbrace{\frac{16\pi G}{c^4} |g| T^{\mu\nu}}^{\text{matter fields}} + \underbrace{\Lambda^{\mu\nu}[h, \partial h, \partial^2 h]}_{\text{non-linearities: } \Lambda \sim h \partial^2 h + \partial h \partial h + h \partial h \partial h + \dots}$$

$$\text{matter eoms: } \nabla_\nu T^{\mu\nu} = 0$$

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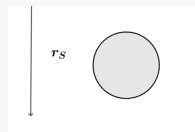
$$\text{matter eoms: } \nabla_\nu T^{\mu\nu} = 0$$

Internal zone $r \leq r_s$

- point-particle approximation

$$S_m = - \sum_a m_a \int d\tau_a$$

- finite-size effects: tides, spins



The b.a-ba of gravitational waves

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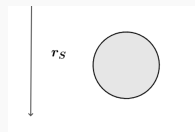
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$$S_m = - \sum_a m_a \int d\tau_a$$

- finite-size effects: tides, spins

Near zone zone $\lambda_{GW} \gg r \gg r_s$

- 2-body dynamics



$$\mathbf{a}_1 = -\frac{G m_2}{r_{12}^3} \mathbf{n}_{12} + O\left(\frac{1}{c^2}\right)$$

The b.a-ba of gravitational waves

Gravitational wave field

$$\nu \equiv \frac{m_1 m_2}{m_1 + m_2}, \quad x = \frac{G(m_1 + m_2)}{rc^2}$$

$$H_{ij}^{TT} = \frac{2G}{c^4 R} P_{ijkl}(\mathbf{N}) \left\{ \ddot{U}_{kl} \left(T - \frac{R}{C} \right) + O\left(\frac{1}{c}\right) \right\}$$

$$U_{ij} = \sum_a m_a x_a^{<ij>}$$

Energy balance equation

$$\begin{aligned} \left\langle \frac{dE}{dt} \right\rangle &= -\langle \mathcal{F} \rangle \quad \text{with} \quad \mathcal{F} \equiv \left(\frac{dE}{dt} \right)^{\text{GW}} = \frac{G}{c^5} \left[\ddot{U}_{ij} \ddot{U}_{ij} + O\left(\frac{1}{c^2}\right) \right] \\ &= \frac{32c^5 \nu^2 x^5}{5G} \end{aligned}$$

The b.a-ba of gravitational waves

Gravitational wave field

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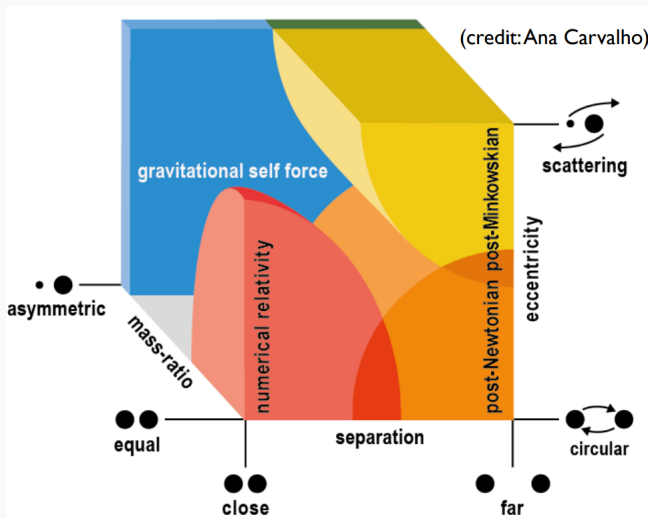
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Dynamics period decay $\frac{dP}{dt}$, eccentricity $\frac{de}{dt}$

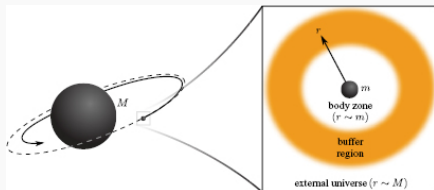
GW modes amplitude $a(t) \propto (t_c - t)^{1/4}$, phase $\phi(t) \propto (t_c - t)^{5/8}$

The different methods

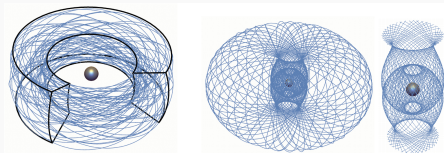


- Inspiral-Merger-Ringdown (IMR): *effective-one-body, phenomenological & surrogate models*

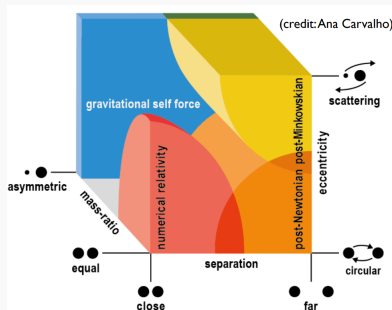
The different methods: gravitational self-force



- extreme mass ratio inspiral
- expansion in $q = \frac{m_1}{m_2} \ll 1$
- resonances, par ex. 2:3



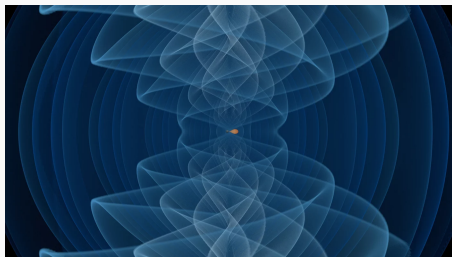
Barack & Pound '18



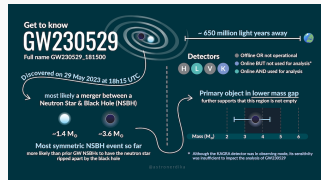
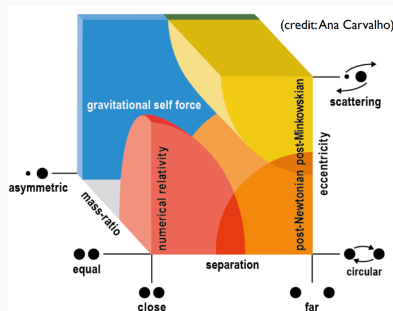
we need to go to **second order self-force** for LISA's accuracy

The different methods: numerical relativity

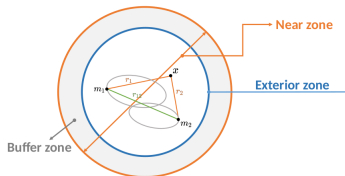
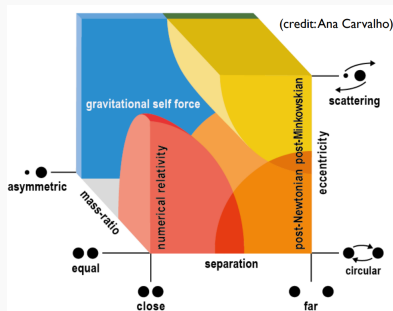
- solving the full Einstein equations
- computationally expensive
- add spins, eccentricity, etc.



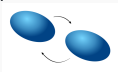
I. Markin, T. Dietrich, H. Pfeiffer, A. Buonanno (Potsdam University and Max Planck Institute for Gravitational Physics)



The different methods: post-Newtonian

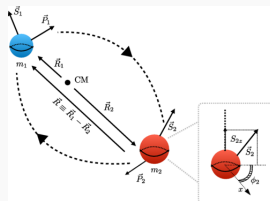


- expansion in $\epsilon = \frac{v_{12}^2}{c^2} \sim \frac{G m_{1,2}}{r_{12} c^2} \ll 1$
- point-particle approximation
- add spins, tides, etc.



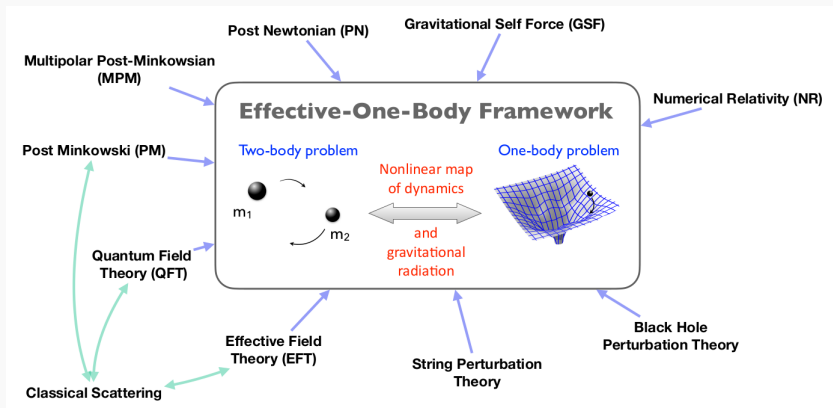
Current state-of-the-art:

- 4.5PN for p.p.**
- NNLO for tides
- $\sim 3\text{PN}$ for spins



Tanay et al. '23

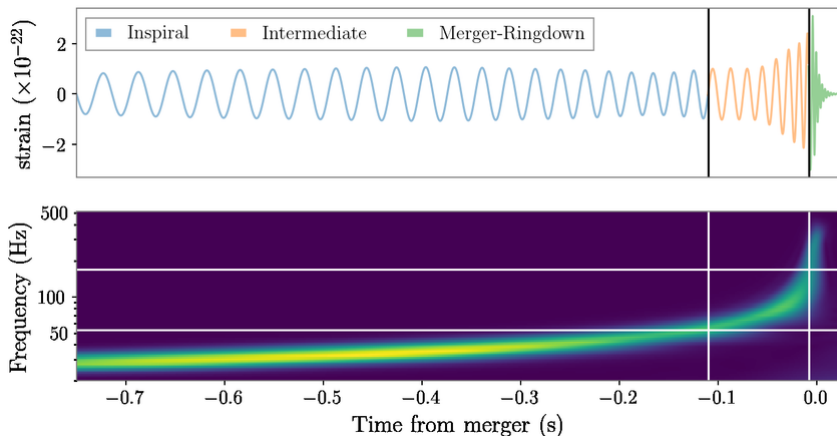
Full IMR waveform: the EOB class



LISA waveform white paper '23

Full IMR waveform: the Phenom class

$$h(f) = \mathcal{A}(f) e^{\psi_n(f)} \quad \psi_n = \{\varphi_{0..7}, \sigma_{0..4}, \beta_{1..3}, \alpha_{0..5}\}$$

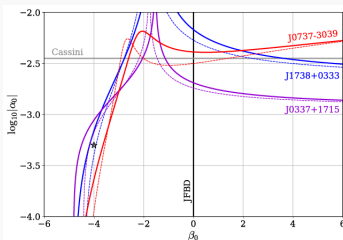


Kwok et al. '21

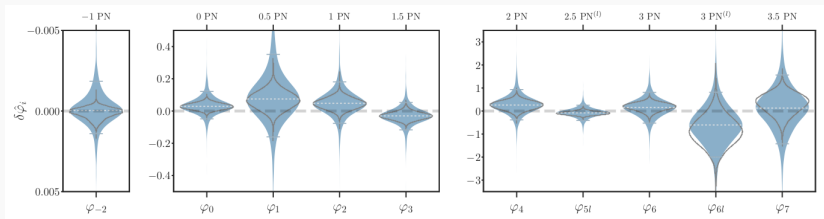
Why do we need...

1. ...to have a bank of **extremely precise** waveform templates?
 - detection and parameter estimation
2. ...to use different modeling techniques
 - different sources, Einstein eqs. hard to solve
3. ...to go **beyond GR**?
 - better understanding of GR
 - to answer fundamental physics and/or cosmology inconsistencies
 - to test the gravitational interaction

GR: a beautiful and successful theory

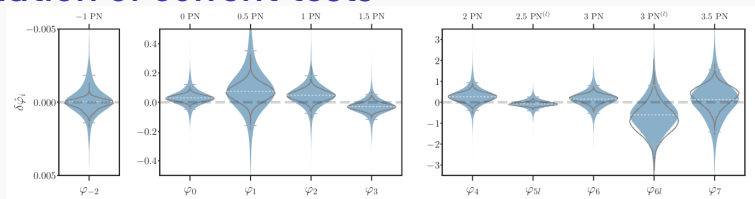


Kramer et al. '21



LIGO-Virgo '21

Limitation of current tests



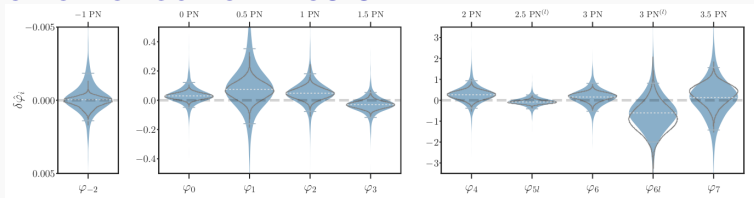
LIGO-Virgo '21

- agnostic searches: $h(f) = \mathcal{A}(f) e^{\psi_n(f) + \delta\psi_n(f)}$, $\delta\psi_n = \{\delta\varphi_{-2,0,\dots,7}\}$
- limited theory-specific tests, ex:
 - ST theories: 3PN dynamics, 2.5PN radiation, NNLO tides
 - EsGB, Chern-Simmons: LO correction, including spin

But:

- ★ we can miss some deviations
- ★ no systematic way to connect deviations to new physics

Limitation of current tests



LIGO-Virgo '21

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What if we had a dictionary?

The dictionary

$$S = \frac{\Lambda^2}{2} \int d^4x \sqrt{-g} \left\{ R + \sum_{p \geq 3} \alpha_{(p-2)} M_\Lambda^{2p-2} (R_{\mu\nu\rho\sigma})^p \right\}$$

- M_Λ new physics scale
- m characteristic mass of the system $\sim$ lightest component

The dictionary

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- M_Λ new physics scale

$$h \sim h_{\text{GR}} \cdot \delta \cdot \left(\frac{v^2}{c^2} \right)^{3p-4, 5^*} \cdot \left(\frac{M_\Lambda}{m} \right)^{2p-2} \cdot \alpha_{(p)}$$

- m characteristic mass of the system $\sim$ lightest component
 - stronger for low mass systems: higher curvature
 - stronger in the inspiral
- leading correction at $3p - 4$ or 5PN (tides) if no extra degree of freedom

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- m characteristic mass of the system $\sim$ lightest component
 - ▶ stronger for **low mass** systems: higher curvature
 - ▶ stronger in the **inspiral**
- leading correction at **$3p - 4$ or 5PN (tides)** if no extra degree of freedom
- ★ systematic way to **distinguish new physics from systematics**
- ★ can be **directly integrated** into existing data analysis methods

Conclusion: we need...

1. ...to have a bank of **extremely precise** waveform templates
 - ▶ detection and parameter estimation
2. ...to use **different modeling techniques**
 - ▶ different sources, Einstein eqs. hard to solve
3. ...to go **beyond GR**
 - ▶ better understanding of gravity, fundamental physics and cosmology questions



Thank you!